### Lagrangian Submanifolds and Constrained Variational Calculus

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"Everything is a Lagrangian Submanifold."

Alan Weinstein.

### Introduction

Lagrangian Submanifolds are very useful!!

#### Definition

#### R. Abraham and J.E. Marsden. "Foundations of Mechanics" (1978)

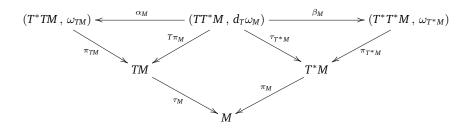
• What is a Lagrangian Submanifold (LS)?

If  $(M, \omega)$  is a symplectic manifold of finite dimension,  $N \subset M$  submanifold,  $i_N : N \hookrightarrow M$ :

- $i_N^* \omega = 0$
- $\dim N = \frac{1}{2} \dim M$
- Examples:
  - $g: M \to M$  is a symplectomorphism, then  $\operatorname{Graph}(g) = \{(x, g(x)), x \in M\} \subset (M \times M, \Omega = pr_1^* \omega pr_0^* \omega) \Rightarrow i_g: \operatorname{Graph}(g) \hookrightarrow M \times M \text{ is a LS.}$
  - $f: M \to \mathbb{R} \Rightarrow df(M) \subset (T^*M, \omega_M)$  is a LS. If (q, p) are local coord. for  $T^*M$ , by Darboux Theorem  $\omega_M = dq \wedge dp$ .
  - $(T^*M, \omega_M)$ ,  $H: T^*M \to \mathbb{R}$ ,  $i_{X_H}\omega_M = dH \Rightarrow X_H(T^*M) \subset (TT^*M, d_T\omega_M)$  is a LS.

### Tulczyjew Triple

• Tulczyjew Triple (1976):



•  $\beta_M = \flat_{\omega_{T^*M}}$ ,  $\beta_{\omega_{T^*M}}(\nu) = i_{\nu}\omega_{T^*M}$ , where  $\nu \in TT^*M \Rightarrow$ 

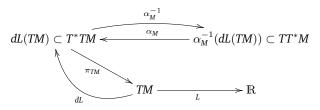
$$\beta_M(q^i, p_i, \dot{q}^i, \dot{p}_i) = (q^i, p_i, -\dot{p}_i, \dot{q}^i).$$

•  $\langle \alpha_M(z), w \rangle = \langle z, \kappa_M(w) \rangle$ , where  $z \in TT^*M$  and  $w \in TTM$ ,

$$\alpha_M(q^i, p_i, \dot{q}^i, \dot{p}_i) = (q^i, \dot{q}^i, \dot{p}_i, p_i).$$

# Intrinsic description of Lag. and Ham. Mechanics

•  $L: TM \to \mathbb{R}$ , then  $dL(TM) \subset T^*TM$  is a LS. Moreover, given that  $\alpha_M$  is a symplectomorphism,  $\alpha_M^{-1}(dL(TM)) \subset TT^*M$  is a LS.



• The solutions of the dynamics of  $\alpha_M^{-1}(dL(TM))$  are curves  $\gamma:I\subset\mathbb{R}\to T^*M$  s.t.  $\frac{d\gamma}{dt}:I\subset\mathbb{R}\to TT^*M$  verifies  $\frac{d\gamma}{dt}(I)\subset\alpha_M^{-1}(dL(TM))$ . Locally

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}}\right) - \frac{\partial L}{\partial q} = 0.$$

### Question

### First Question

Is there a way to describe constrained Lagrangian Mechanics through Lagrangian Submanifolds and, moreover, to relate that dynamics with a Hamiltonian system?

# Lagrangian Mechanics with Constraints

• Constrained Lagrangian System (CLS):  $C \subset TM$  and a Lag. function  $L: C \to \mathbb{R}$ . Locally C is described by a set of constraints  $\phi^{\alpha}: TM \to \mathbb{R}$ .

### Theorem (Tulczyjew)

Let M be a smooth manifold,  $N \subset M$  a submanifold, and  $f: N \to \mathbb{R}$  a smooth function. Then

$$\Sigma_f = \left\{ p \in T^*M \mid \pi_M(p) \in N \text{ and } \langle p, \nu \rangle = \langle \mathrm{d}f, \nu \rangle \right.$$
 for all  $\nu \in TN \subset TM$  such that  $\tau_M(\nu) = \pi_M(p) \right\}$ 

is a Lagrangian submanifold of  $T^*M$ .

• Hence,  $\boxed{\Sigma_L \in (T^*TM\,,\,\omega_{TM})}$  is a LS. Moreover,  $\boxed{\alpha_M^{-1}(\Sigma_L) \subset (TT^*M\,,\,d_T\omega_M)}$  also is, which means DYNAMICS!. Locally  $\mathbb{L} = \tilde{L} + \lambda_\alpha \phi^\alpha : TM \to \mathbb{R}$ ,

$$\frac{d}{dt} \left( \frac{\partial \tilde{L}}{\partial \dot{q}^{i}} + \lambda_{\alpha} \frac{\partial \phi^{\alpha}}{\partial \dot{q}^{i}} \right) - \frac{\partial \tilde{L}}{\partial q^{i}} - \lambda_{\alpha} \frac{\partial \phi^{\alpha}}{\partial q^{i}} = 0$$

$$\phi^{\alpha} (q^{i}, \dot{q}^{i}) = 0$$

# From Constrained Lagrangian Mechanics to Hamiltonian Mechanics

### **Constrained Legendre Transformation**

 $\mathbb{F}L: \Sigma_L \longrightarrow T^*M$  as the mapping  $\mathbb{F}L = \tau_{T^*M} \circ (\alpha_M^{-1})|_{\Sigma_L}$ .

We will say that (L,C) is regular if  $\mathbb{F}L$  is a local diffeomorphism and hyperregular if  $\mathbb{F}L$  is a global diffeomorphism.

- $E_L: \Sigma_L \to \mathbb{R}$ ,  $E_L(\alpha_u) = \langle \alpha_u, u_u^V \rangle L(u)$ , where  $\alpha_u \in \Sigma_L$ ,  $u \in C$  and  $u_u^V \in TC \subset TTM$ .
- Define  $\omega_L = (\mathbb{F}L)^* \omega_M$  on  $\Sigma_L \Rightarrow i_X \omega_L = dE_L$ .
- ullet If (L,C) is hyperrregular, we can define a Hamiltonian function  $H:T^*M \to \mathbb{R}$  by

$$H = E_L \circ (\mathbb{F}L)^{-1}$$

such that  $i_{X_H}\omega_M=dH$ .

$$\Sigma_L \Rightarrow X_H(T^*M)$$

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$$\alpha_M^{-1}(\Sigma_L) = X_H(T^*M)$$

# From Hamiltonian Mechanics to Constrained Lagrangian Mechanics

•  $H: T^*M \to \mathbb{R}$ . Since  $\pi_M: T^*M \to M$  is a vector bundle, define the dilation vector field  $\Delta^* \in \mathfrak{X}(T^*M)$ :

$$\Delta^*(\alpha_q)f = \frac{d}{dt}\Big|_{t=0} f(t\alpha_q), \quad \Delta^* = p_i \frac{\partial}{\partial p_i}.$$

• Define the fiber derivative  $\mathbb{F}H: T^*M \to TM$ 

$$\langle \mathbb{F}H(\alpha_q), \beta_q \rangle = \frac{d}{dt}\Big|_{t=0} H(\alpha_q + t\beta_q), \quad \mathbb{F}H(q^i, p_i) = (q^i, \frac{\partial H}{\partial p_i}).$$

• Assume  $\mathbb{F}H(T^*M)=C$  is a submanifold  $\subset TM\Rightarrow$  Mimicking the Gotay and Nester's definition, we set  $L:C\to\mathbb{R}$  such that

$$L \circ \mathbb{F}H = \Delta^*H - H.$$

•  $\mathbb{F}H(\alpha_q) = \mathbb{F}H(\beta_q)$  or equivalently  $(\Delta^*H - H)(\alpha_q) = (\Delta^*H - H)(\beta_q)$ . Extra assumption

#### Definition

A Hamiltonian  $H: T^*M \to \mathbb{R}$  is **almost-regular** if  $\mathbb{F}H(T^*M) = C$  is a submanifold of TM and  $\mathbb{F}H: T^*M \to C \subset TM$  is a submersion with connected fibers.

# From Hamiltonian Mechanics to Constrained Lagrangian Mechanics

• Under this assumption we can consider the infinitesimal condition, i.e.

$$\mathcal{L}_Z(\Delta^*H - H) = 0, \ \forall \ Z \in \ker(\mathbb{F}H_*).$$

• This is the case, hence *L* is well defined.

#### Theorem

The following equivalence holds

$$\alpha_M(X_H(T^*M)) = \Sigma_L .$$

#### Sketch of the Poof

For any  $W_1 \in X_H(T^*M)$ 

$$\langle \alpha_M(W_1), U \rangle = \langle dL, U \rangle \ \forall U \in TC.$$

This is equivalent to prove

$$\langle \alpha_M(W_1), \mathbb{F}H_*(W_2) \rangle = \langle dL, \mathbb{F}H_*(W_2) \rangle \ \forall W_2 \in TT^*M.$$

$$\mathbb{F}H^*\alpha_M(W_1) = \mathbb{F}H^* dL = d(\Delta^*H - H).$$

### From Hamiltonian Mechanics to Constrained Lagrangian Mechanics

$$TT^*M \supset X_H(T^*M) \xrightarrow{\alpha_M} \alpha_M(X_H(T^*M)) = \Sigma_L \hookrightarrow T^*TM$$

$$\downarrow^{X_H} \uparrow \qquad \qquad \downarrow^{dL}$$

$$T^*M \qquad \qquad C \subset TM$$

$$\downarrow^L \qquad \qquad \mathbb{R}$$

# ANY DISCRETE EQUIVALENCE?

### **Brief Introduction to Discrete Mechanics**

#### J.E. Marsden and M. West: "Discrete Mechanics and Variational Integrators" (2001)

• *M* is a *n*-dimensinal smooth manifold  $(q^i)$ , and *TM* its tangent bundle  $(q^i, \dot{q}^i)$ .

 $L:TM\to\mathbb{R}$ :

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}^i}\right) - \frac{\partial L}{\partial q^i} = 0, \quad 1 \le i \le n.$$

• Discrete Mechanics:  $TM \Rightarrow M \times M$ . Moreover  $L_d : M \times M \to \mathbb{R}$ :

$$L_d(q_0,q_1) pprox \int_0^h L(q(t),\dot{q}(t)) dt$$

• Define the sequences  $(q_0,...,q_N) \in M^{N+1}$  and the action sum  $S_d = \sum_{k=1}^N L_d(q_{k-1},q_k)$ .

$$D_1L_d(q_k, q_{k+1}) + D_2L_d(q_{k-1}, q_k) = 0.$$

This equations define a discrete flow  $\phi_{L_d}: M \times M \to M \times M$ ,  $\phi_{L_d}(q_{k-1},q_k) = (q_k,q_{k+1})$ , under regularity assumptions.

#### Brief Introduction to Discrete Mechanics

• Go into a Hamiltonian Picture through Discrete Legendre Transforms

- ullet Pullback the canonical 2-form:  $\omega_d=(\mathbb{F}L_d^-)^*\omega_M=(\mathbb{F}L_d^+)^*\omega_M$
- Preservation properties
  - $\phi_{L_d}^* \omega_d = \omega_d \Rightarrow$  Symplectic Preservation.
  - The flow  $\phi_{L_d}$  also preserves  $J_d: M \times M \to \mathfrak{g}^*$ , defined by

$$\langle J_d(q_k, q_{k+1}), \xi \rangle = \langle D_2 L_d(q_k, q_{k+1}), \xi_M(q_{k+1}) \rangle$$

where  $\xi_M$  is the fundamental vector field determined by  $\xi \in \mathfrak{g} \Rightarrow$  Momentum Map Preservation.

### Symplectic Hamiltonian Numerical Methods

•  $H: T^*M \to \mathbb{R} \Rightarrow \text{HAMILTON EQUATIONS} \Rightarrow \phi_t: T^*M \to T^*M, \ (q(0), p(0)) \in T^*M$ 

$$\phi_t(q(0), p(0)) = (q(t), p(t))$$

• What if the equations are not easily solvable (or not solvable at all)?  $\Rightarrow$  NUMERICAL METHOD  $\Rightarrow \phi_h : T^*M \to T^*M$ ,  $(q_0, p_0) = (q(0), p(0)) \in T^*M$ 

$$\phi_h(q_0, p_0) = (q_1, p_1),$$

•  $(q_1, p_1) \simeq (q(h), p(h))$ . We can manage to obtain a symplectic method, i.e.

$$\omega_0 = \omega_1$$
, that is  $dq_0 \wedge dp_0 = dq_1 \wedge dp_1$ .

• For instance: Symplectic Euler methods

$$q_1 = q_0 + h H_p(q_0, p_1), \quad p_1 = p_0 - h H_q(q_0, p_1)$$

Hairer, Lubich, Wanner: "Geometric Numerical Integration: Structure Preserving algorithms for ODEs" (2002)

# Symplectic Hamiltonian Numerical Methods

$$\phi_h: T^*M \to T^*M$$

$$((q_0,p_0),(q_1,p_1))\in T_{q_0}^*M\times T_{q_1}^*M$$

J.C. Marrero, D. Martín de Diego and A. Stern. "Symplectic groupoids and discrete constrained Lagrangian mechanics" (2011).

- Discrete Constrained Lagrangian (DCL) system can be defined  $(C_d, L_d)$ ,  $C_d \subset M \times M$  is a submanifold,  $L_d : C_d \to \mathbb{R}$  is the discrete Lagrangian function.
  - Thus  $\Sigma_{L_d} \subset (T^*(M \times M) \, , \, \omega_{M \times M})$  is a LS by Tulczyjew's Theorem.
- The proper scenario to describe Hamiltonian symplectic numerical methods is  $(T^*M \times T^*M\,,\,\Omega)$ , where  $\Omega = pr_1^*\,\omega_M pr_0^*\,\omega_M$ , and  $pr_i: T^*M \times T^*M \to T^*M$

### Second Question

Is there a way to relate a discrete (constrained) Lagrangian Submanifold and a general discrete symplectic Hamiltonian dynamics?

$$\Lambda \subset T^*M \times T^*M \sim \Sigma_{L_d}??$$

Symplectomorphism

$$\Upsilon: \left(T^*(M \times M), \, \omega_{M \times M}\right) \quad \to \quad \left(T^*M \times T^*M, \, \Omega\right) \\ \gamma_{(q_0, q_1)} \equiv \left(\gamma_{q_0}, \gamma_{q_1}\right) \quad \mapsto \quad \left(-\gamma_{q_0}, \gamma_{q_1}\right)$$

where  $(q_0, q_1) \in M \times M$  and  $\gamma_{q_i} \in T^*M$ .

- We generate  $\Upsilon(\Sigma_{L_d})$  LS of  $T^*M \times T^*M$ .
- Dynamics:  $\gamma_{q_0},...,\gamma_{q_N}$  s.t.  $(\gamma_{q_i},\gamma_{q_{i+1}})\in \Upsilon(\Sigma_{L_d}), 0\leq i\leq N-1$ .

$$\gamma_{q_k} \in T_{q_k}^* M \cap pr_0(\Upsilon(\Sigma_{L_d})) \cap pr_1(\Upsilon(\Sigma_{L_d})), 1 \le k \le N-1.$$

D. Iglesias, J.C. Marrero, D. Martín de Diego and E. Padrón: "Discrete Dynamics in Implicit Form" (2011).

### Constrained Discrete Legendre Transformations

The mappings  $\mathbb{F}L_d^{\pm}: \Sigma_{L_d} \longrightarrow T^*M$  are defined by

$$\begin{array}{rcl} \mathbb{F}L_d^- & = & pr_0 \circ \Upsilon\big|_{\Sigma_{L_d}}, \\ \mathbb{F}L_d^+ & = & pr_1 \circ \Upsilon\big|_{\Sigma_{L_d}}. \end{array}$$

We'll say that  $(L_d, C_d)$  is regular if  $\mathbb{F}L_d^-$  is a local diffeomorphism and hyperregular if  $\mathbb{F}L_d^-$  is a global diffeomorphism.

• Generating function: if N LS  $N\subset (M,\omega=d\theta)$ , where  $\theta$  is the Liouville 1-form. Then holds that

$$0=i_N^*\omega=d(i_N^*\theta),$$

and consequently, by Poincaré's Lemma  $i_N^*\theta = dS$ . S is the generating function of N.

- Consider  $\Lambda$  LS  $\Lambda \subset T^*M \times T^*M$ .
- Consider  $C_d \subset M \times M$ ,  $L_d : C_d \to \mathbb{R}$ , and the LS  $\Sigma_{L_d}$ .
- Consider  $\Upsilon^{-1}(\Lambda) \subset (T^*(M \times M), \omega_{M \times M}).$
- Assume  $i_{\Upsilon^{-1}(\Lambda)}^* \theta_{M \times M} = dS$ , i.e. there exists a generating function S of  $\Upsilon^{-1}(\Lambda)$ .
- Assume that  $\pi_{M\times M}(\Upsilon^{-1}(\Lambda))=C_d$  is a submersion with connected fibers.

#### Theorem

Under the previous conditions the function  $S: \Upsilon^{-1}(\Lambda) \to \mathbb{R}$  is  $(\pi_{M \times M})\big|_{\Upsilon^{-1}(\Lambda)}$ -projectable onto a function  $L_d: C_d \to \mathbb{R}$ . Moreover, the following equation holds

$$\Upsilon^{-1}(\Lambda) = \Sigma_{L_d} .$$

 $\Lambda \subset T^*M \times T^*M \sim \Sigma_{L_d} \text{ YES!!}$ 

Example: Martinet type sub-Riemannian structure (continuous case)

$$H(q,p) = \frac{1}{2} \left( \left( p_x + p_z \frac{y^2}{2} \right)^2 + \frac{p_y^2}{(1+\beta x)^2} \right),$$

- $q = (x, y.z)^T \in \mathbb{R}^3$  and  $p = (p_x, p_y, p_z) \in (\mathbb{R}^3)^* \simeq \mathbb{R}^3$ .
- $\mathbb{F}H(x, y, z; p_x, p_y, p_z) = (x, y, z; (p_x + p_y \frac{y^2}{2}), \frac{p_y}{(1+\beta x)^2}, (p_x + p_z \frac{y^2}{2}) \frac{y^2}{2}) = (x, y, z; \dot{x}, \dot{y}, \dot{z})$   $C \subset T\mathbb{R}^3 = \left\{ (x, y, z; \dot{x}, \dot{y}, \dot{z}) \text{ s.t. } \dot{z} = \frac{y^2}{2} \dot{x} \right\}.$
- $\bullet L \circ \mathbb{F}H = \Delta^*H H$ :

$$L(q, \dot{q}) = \frac{1}{2} \left( \dot{x}^2 + (1 + \beta x)^2 \dot{y}^2 \right).$$

• Corresponds to the sub-Riemannian structure  $(\Delta, g)$ , being  $\Delta = \ker \alpha$  for  $\alpha = dz - \frac{y^2}{2} dx$  and  $g = dx^2 + (1 + \beta x)^2 dy^2$ . It is clear that  $L(q, \dot{q}) = \frac{1}{2} g(\partial/\partial q, \partial/\partial q)$ .

### Example: Martinet type sub-Riemannian structure (discrete case)

• We apply a symplectic Euler method to the Martinet Hamiltonian

$$q_1 = q_0 + h H_p(q_0, p_1), \quad p_1 = p_0 - h H_q(q_0, p_1).$$

• We define  $\Upsilon^{-1}(\Lambda)$  by means of the generating function  $H^+(q_0,p_1)=q_0\,p_1+hH(q_0,p_1)$  i.e.

$$\begin{array}{rcl} x_1 & = & x_0 + h\left((p_1)_x + (p_1)_z\frac{y_0^2}{2}\right)\right),\\ \\ y_1 & = & y_0 + h\frac{(p_1)_y}{(1+\beta x_0)^2},\\ \\ z_1 & = & z_0 + h\left((p_1)_x + (p_1)_z\frac{y_0^2}{2}\right)\right)\frac{y_0^2}{2}, \end{array}$$

- which determines  $C_d \subset M \times M \Rightarrow (z_1 z_0) = \frac{y_0^2}{2}(x_1 x_0)$ .
- Finally, we find  $S(q_0, q_1) = h\left(p_1 \frac{\partial H(q_0, p_1)}{\partial p_1} H(q_0, p_1)\right)$  which is projectable onto

$$L_d(q_0,q_1) = \frac{1}{2h} \left( \left( x_1 - x_0 \right)^2 + \frac{y_0^2}{2} \left( y_1 - y_0 \right)^2 \right) = h L \left( q_0, \frac{q_0 - q_1}{h} \right).$$

#### Conclusions

We have shown that given a CLS one can always find a Hamiltonian function.
 Moreover, we prove that given an arbitrary Hamiltonian system one can always construct a (possibly) CLS that generates the original system.

• We try to get some light over discrete mechanics, which can be interpreted as suitable LS of  $T^*M \times T^*M$ . We geometrically find when the discrete variational procedure ( $\Sigma_{L_d}$ ) matchs a symplectic numerical method for the associated Hamiltonian system ( $\Lambda$ ).

# THANKS!!